## The Surface Structure of the Active Bismuth Molybdate Catalyst

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 $Bi_2Mo_3O_{12}(\alpha)$ ,  $Bi_2Mo_2O_9(\beta)$ ,  $Bi_2MoO_6(\gamma)$ , and mixtures of  $\alpha$  and  $\gamma$  were investigated as to their activity for the oxidative dehydrogenation of but-1-ene to butadiene. Pure  $\gamma$  was found to be almost inactive,  $\alpha$  was only moderately active, but  $\beta$  and all mixtures of  $\alpha$  and  $\gamma$  showed considerable and approximately equal activity. XPS measurements showed the surface of  $\alpha$ ,  $\beta$ , and  $\gamma$  to have similar Bi/Mo ratios to the bulk. However, for mixtures of  $\alpha$  and  $\gamma$  with average ratios 2/3 < Bi/Mo < 2/1, the Bi/Mo ratio at the surface was found to be approximately equal to 1, decreasing steeply at the Mo-rich side and increasing even more steeply at the Bi-rich edge. The catalytic activity therefore seems to be connected with a surface layer with a Bi/Mo ratio near 1. Adsorption studies on inactive and active catalysts near the Bi/Mo = 2 edge show that B-type adsorption (weak olefin adsorption) is already found on inactive catalyst (stoichiometric y). A-Type adsorption (strong butadiene adsorption) only becomes observable when a considerable amount of excess MoO<sub>3</sub> is introduced; the amounts adsorbed increase almost linearly with the increase in activity of the catalyst. In infrared and Raman spectroscopic studies, the theoretical positions for the bands related to Mo-O vibrations were determined, using atomic coordinates from the literature and stretching force constants obtained by the application of the Cotton-Wing relation and finally computing the position via a modified Schachtschneider program. The information thus acquired for compounds with known structure was used to analyze the spectra of Bi-molybdates with unknown structure. The results indicate that the active nonstoichiometric catalysts are characterized by the presence of a surface layer on  $\gamma$ . The new bands formed in catalysts containing small amounts of excess MoO<sub>3</sub> were found to be similar to those of  $\beta$ . It is concluded that the surface layer has a structure closely related to that of  $\beta$ . A model of the surface of an active catalyst is given.

#### INTRODUCTION

The catalytic activity of Bi-molybdates in the oxidation of olefins has been studied by many investigators (1-4). Only binary oxides with an atomic ratio in the range 2/3< Bi/Mo < 2/1 were found to be active while outside this range there is hardly any activity. The Bi<sub>2</sub>Mo<sub>3</sub>O<sub>12</sub> ( $\alpha$ ), Bi<sub>2</sub>Mo<sub>2</sub>O<sub>9</sub> ( $\beta$ ), and Bi<sub>2</sub>MoO<sub>6</sub> ( $\gamma$ ) are present in catalysts within this range of composition. Opinions vary as to which compound is the active catalyst. Kolchin *et al.* (5) observed  $\alpha$  and  $\beta$  to have appreciable activity for both oxidation and ammoxidation of propene,

<sup>2</sup> Present address: Asahi Chemical Industry Co. Ltd., Tokyo, Japan. while  $\gamma$  displayed only low activity and poor selectivity. On the other hand, Batist *et al.* (6) reported  $\gamma$  to be comparable in activity to  $\beta$  for the oxidation of butene to butadiene, while  $\alpha$  was only moderately active. Since the three reactions are considered to proceed by similar mechanisms, the conflicting data make it difficult to connect the structure of active sites at the surface unambiguously with a special bulk structure.

The present contribution is an attempt to obtain a better insight into this problem by the application of XPS, ir, and Raman spectroscopy.

#### **EXPERIMENTAL**

Catalyst preparation. Bi<sub>2</sub>MoO<sub>6</sub> ( $\gamma$ ) and Bi<sub>2</sub>Mo<sub>3</sub>O<sub>12</sub> ( $\alpha$ ) were prepared according to a method proposed by Batist *et al.* (6). Solid basic bismuth nitrate BiO(NO<sub>3</sub>) and molyb-

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dic acid  $H_2MoO_4$  were powdered, slurried in water, and boiled under stirring for 16 hr. Thereafter, the solid materials were filtered and dried. The initial ratio of Bi to Mo was 2:1 for the synthesis of  $\gamma$  and 1:5 for  $\alpha$ . In order to remove the excess of MoO<sub>3</sub> for the latter case, the calcined sample was washed with a diluted ammonia solution.  $Bi_2Mo_2O_9$  $(\beta)$  was prepared by coprecipitation from bismuth nitrate  $Bi(NO_3)_3 \cdot 5H_2O$  and ammonium hepta-molybdate  $(NH_4)_6Mo_7O_{27}$  ·  $4H_2O$ , as described by Batist *et al.* (6). The samples of the various oxides were calcined at 500°C for 4 hr. The oxides will be referred to as compounds  $\alpha$ ,  $\beta$ , and  $\gamma$  as a convenient short notation.

Catalysts with other atomic ratios of Bi and Mo,  $0.58 \leq \text{Bi/Mo} \leq 2.04$ , were also prepared, using the same method as for the compound  $\gamma$ . In this paper, they will be called "multiphase" samples, although Xray analysis sometimes showed only the presence of one compound. The atomic ratio of Bi to Mo in these samples was determined by means of atomic absorption spectrometric analysis.

We also prepared some catalysts by impregnating stoichiometric  $Bi_2MoO_6(\gamma)$  with  $MoO_3$  or  $Bi_2O_3$  ("impregnated" samples). Previously prepared  $\gamma$  samples were slurried in a solution of  $H_2MoO_4$  or  $Bi(NO_3)_3 \cdot$  $5H_2O$  of the appropriate concentration. After evaporation under stirring, the solids were calcined at 520°C for 2 hr.

X-Ray analysis showed only the expected phase in the  $\alpha$  and  $\gamma$  compounds. For compound  $\beta$ , small amounts of  $\alpha$  and  $\gamma$  were present. In the multiphase samples, only a mixture of compounds  $\alpha$  and  $\gamma$  were found, without compound  $\beta$  being observable. Samples near 2/1, with compositions situated between an excess of 4% Bi to an excess of 8% Mo, showed only the lines of the  $\gamma$ -phase.

The surface areas of the various samples are plotted versus the atomic ratio of Bi/Mo in Fig. 1. Densities also shown in this figure are determined pycnometrically with carbon tetrachloride. The values of the densities of the three compounds are close to the densities calculated from the unit cell dimension of the crystal structure.

	Obs	Calcd
$Bi_2MoO_6(\gamma)$	7.8	8.0 (7)
$Bi_2Mo_2O_9(\beta)$	6.5	6.6 (8)
$Bi_2Mo_3O_{12}(\alpha)$	6.0	6.2 (9)

It is noteworthy that multiphase samples usually had higher surface areas and lower densities than the pure compounds.

Catalyst activity measurements. The catalytic activity of the various samples for the oxidation of but-1-ene to butadiene was determined for a range of temperatures from 350°C upwards. The apparatus, procedure, gas chromatographic analysis, and the method of calculation were given earlier by Batist *et al.* (10). In each experiment sample weights of unit surface area were used, the gas flows were 20 cm<sup>3</sup> min<sup>-1</sup> of butene and 100 cm<sup>3</sup> min<sup>-1</sup> of artificial air.

But-1-ene and butadiene adsorption measurements. A conventional vacuum apparatus equipped with a quartz adsorption cell was used for isotherm measurements. The measurements were performed over a range of pressures from  $10^{-3}$  to 10 Torr and a range of temperatures from 0 to 200°C. The adsorption method was described elsewhere (11).

XPS measurements. X-Ray photoelectron spectroscopy was performed using an



FIG. 1. Surface areas and bulk densities of various "multiphase" samples and "compounds" of Bi-molybdates as a function of composition.

AEI ES 200 spectrometer with MgK $\alpha$ . The power supply was run at 12 kV and 15 mA, the spectra being collected with a PPDS computer and stored in a 32 K disc. The powdered samples were mounted on adhesive tape.

Infrared and raman spectra measurements. The ir spectra of samples were recorded with a Grubb Parsons spectrometer MK-3 in the range of 1250-400 cm<sup>-1</sup> and a Hitachi model EP 1-92 in the range of 700-200 cm<sup>-1</sup>, using the conventional KBr disk techniques. The Raman spectra of powders from 1250 to 0 cm<sup>-1</sup> were obtained using a Varian carey-82 with an Ar laser (5145 Å). Samples were either measured on a rotating disk or in glass capillary tubes; no essential difference was observed between spectra observed by the two methods.

### EXPERIMENTAL RESULTS

Catalytic activity measurements. Figure 2 shows the activities of the catalysts. For all catalysts, the selectivity to butadiene was found to be higher than 95% over the whole temperature range. The compound  $\gamma$  was found to be only weakly active. Multiphase samples around 2/1 with a small excess of Mo are fully active, while those with a small excess of Bi are entirely inactive. The decline at the 2/1 edge appears to be very sharp. The compound  $\alpha$  is only moderately active, but a slight excess of Bi brings the sample to an activity level com-



FIG. 2. Conversion of but-1-ene (per unit surface area) as a function of composition for "multiphase" samples and "compounds" of Bi-molybdates.



FIG. 3. Conversion of but-1-ene (per unit surface area at 440°C) as a function of addition of  $MoO_3$  or  $Bi_2O_3$  on stoichiometric  $Bi_2MoO_6$  (impregnated samples).

parable to that of the compound  $\beta$  and all samples with  $0.8 \leq \text{Bi/Mo} \leq 1.92$  show similar activities at temperatures higher than 440°C. However, at temperatures lower than 410°C, the activities decrease with decreasing of Bi/Mo ratio.

The activities of the impregnated samples at 440°C are shown in Fig. 3. The results show equally spectacular changes as found with the multiphase samples in the corresponding range. Addition of 4% MoO<sub>3</sub> (atomic percent with respect to Mo, i.e.,  $Bi_2MoO_6 \rightarrow Bi_2Mo_{1.04}O_{6.12}$ ) causes an activity increase by a factor of at least 10. Since, in impregnation, the excess MoO<sub>3</sub> may be assumed to be deposited at the surface, it might be surmised that the multiphase samples also have an excess Mo at the surface.

Adsorption measurements. The adsorption of but-1-ene over all "impregnated" samples is found to be always weak, fast, and reversible. A small excess Bi in the 2/1 surface shows an associative chemisorption with a rather weak heat of adsorption at 25 kJ mole<sup>-1</sup>, similar to that reported earlier by one of us on Bi<sub>2</sub>O<sub>3</sub> (11). Stoichiometric Bi<sub>2</sub>MoO<sub>6</sub> ( $\gamma$ ) has a different type of adsorption with a higher heat of adsorption at 46 kJ mole<sup>-1</sup> and the characteristics of a dissociative adsorption. This adsorption is closely similar to that described earlier by one of us as B-type adsorption (11). Addition of MoO<sub>3</sub> shows no change of the ad-

Catalyst	Composition	Type of isotherm	V <sub>m</sub> (cm <sup>3</sup> m <sup>-2</sup> )	Log P <sub>o</sub> <sup>0</sup> (mm Hg)	Heat of adsorption Q (kJ mol <sup>-1</sup> )
Bi <sub>2</sub> MoO <sub>6</sub>	C <sub>4</sub> H <sub>8</sub>	DS(w) <sup>a</sup>	0.141	8.9	42.2
-	$C_4H_6$	DS(W)	0.138	6.6	41.9
		Strong ads.	no obs.		
+ 1% $Bi_2O_3$	C₄H <sub>8</sub>	$SS(w)^b$	0.106	4.7	25.2
	C <sub>4</sub> H <sub>6</sub>	SS(w)	0.108	4.7	23.1
		Strong ads.	no obs.		
+ 1% MoO <sub>3</sub>	$C_4H_8$	DS(w)	0.133	9.2	46.2
	$C_4H_6$	DS(w)	0.135	6.6	42.1
		Strong ads.	0.013		
+ 2% MoO <sub>3</sub>	$C_4H_8$	DS(w)	0.130	9.5	46.2
	C <sub>4</sub> H <sub>6</sub>	DS(w)	0.130	6.9	42.0
	• •	Strong ads.	0.023		
+ 4% MoO <sub>3</sub>	C <sub>4</sub> H <sub>8</sub>	DS(w)	0.116	9.5	46.2
	C <sub>4</sub> H <sub>6</sub>	DS(w)	0.118	6.9	41.6
		Strong ads.	0.032		

TABLE 1

Adsorption Data of But-1-ene and Butadiene on Impregnated Samples

<sup>a</sup> DS, dual-site adsorption.

<sup>b</sup> SS, single-site adsorption.

sorption parameters but decreases the surface concentration. The activated and strong desorption of butadiene, A-type adsorption, described in the earlier paper (11), is not present on Bi-enriched or stoichiometric  $Bi_2MoO_6$  ( $\gamma$ ). It only becomes observable when a considerable amount of excess  $MoO_3$  is introduced. The adsorption parameters for the Langmuir-Hinshelwood isotherm are given in Table 1.

XPS measurements. Figure 4 shows the XPS spectra of compounds  $\alpha$ ,  $\beta$ , and  $\gamma$  and of two multiphase samples near the 2/1 edge but with a slight excess of Mo or Bi. To calculate the Bi/Mo atomic ratio at the surface we used the method proposed by Carter *et al.* (12) who applied the photoelectron cross sections of Scofield (13) for AlK $\alpha$  or MgK $\alpha$  radiation. The relative intensity N for two different photoelectron peaks from the same samples is given as

$$N_1/N_2 = \sigma_1 n_1 \lambda_1 S_1/\sigma_2 n_2 \lambda_2 S_2, \qquad (1)$$

where  $\sigma$  is the photoelectron cross section for a given subshell of a given element, *n* is the concentration of atoms per unit volume, S is the spectrometer factor, and  $\lambda$  is the mean free path. Assuming that the ratio of the S factor is not dependent on the energy



FIG. 4. XPS spectra of "compound" Bi-molybdates and "multiphase" samples near the 2/1 edge.

and that the mean free paths can be approximately fitted by the equation  $\lambda \propto E^{1/2}$ , we arrive (for  $n_1 = n_2$ ) at the following relation:

$$N_1/N_2 = \sigma_1(h\nu - E_{B1})^{1/2}/\sigma_2(h\nu - E_{B2})^{1/2}, \quad (2)$$

where  $E_B$  is the energy of the level under investigation and  $h\nu$  is the energy of the radiation quanta.

The value of the relative intensity ratio of  $N_{\rm Bi}/N_{\rm M0}$  is found to be 2.5 from Eq. (2), taking the value of  $\sigma_{\rm Bi,4\,f(7/2)}$  and  $\sigma_{\rm M0,3\,d(5/2)}$  to be respectively 14.0 and 5.77 as calculated by Scofield (13), and  $E_{B,\rm Bi,4\,f(7/2)}$  and  $E_{B,\rm M0,3\,d(5/2)}$  as equal to respectively 159 and 232 eV.

Comparing the Bi/Mo ratios thus predicted with the experimental values for the compounds  $\alpha$ ,  $\beta$ , and  $\gamma$ , there is excellent agreement between bulk compositions and surface compositions, as shown in Fig. 5. However, for all multiphase samples the Bi/Mo ratio in the catalytically active range is found to be near 1, similar to the value of the compound  $\beta$  (see Fig. 5).

Infrared and raman spectra. The ir bands of compound  $\gamma$  and 2/1 with excess MoO<sub>3</sub> are shown in Fig. 6a. The band positions of the stoichiometric compound are similar to those given in the literature (14). Above 600 cm<sup>-1</sup>, the "excess MoO<sub>3</sub>" samples exhibit a new set of weak bands superimposed upon the existing ones. There is a sharp, weak band at 832 cm<sup>-1</sup>, while the main 735-



FIG. 5. Atomic ratio of Bi/Mo at surface from XPS analysis as a function of bulk composition for "multiphase" samples and "compound" Bi-molybdates.

 $cm^{-1}$  band broadens considerably with a shoulder at 760  $cm^{-1}$  as shown in Fig. 7.

The Raman spectrum of compound  $\gamma$  (Fig. 6a) shows three bands at 848, 799, and 715 cm<sup>-1</sup>, of which the second is strong. Between 200 and 400 cm<sup>-1</sup>, there are two moderately strong bands at 325 and 272 cm<sup>-1</sup> with some smaller bands. The range of 400 to 600 cm<sup>-1</sup> is practically empty. Excess MoO<sub>3</sub> causes the appearance of new sharp bands at 886 cm<sup>-1</sup> and two weak shoulders at 809 and 782 cm<sup>-1</sup>. Figure 7 shows increasing intensity of these new bands with addition of MoO<sub>3</sub> in Bi<sub>2</sub>MoO<sub>6</sub> ( $\gamma$ ).

Compound  $\gamma$  heated at 645°C (X-ray diagram shows the Blasse phase,  $\gamma'$ ) are more complicated than those of koechlinite ( $\gamma$ ) over the whole range: there are at least nine bands in the ir and Raman spectra above 600 cm<sup>-1</sup> (see Fig. 6b).

The spectra of compound  $\beta$  are shown in Fig. 6c. The ir spectra show a sharp band at 890 cm<sup>-1</sup> and a much stronger and broader band between 650 and 850 cm<sup>-1</sup>; it is also possible to observe shoulders (840, 710 cm<sup>-1</sup>) and peaks (800, 770, 740 cm<sup>-1</sup>). The Raman spectra show the presence of two strong bands (892, 793 cm<sup>-1</sup>) with two weaker bands at 845 and 715 cm<sup>-1</sup>.

The band positions of the complicated spectrum for compound  $\alpha$  are similar to those in the literature (15). The positions of ir and Raman bands observed for the compounds are summarized in Table 2.

A short survey of the results for other compounds follows below (band positions in  $cm^{-1}$ ):

- $MoO_3$ : 990(s), 875(s), 815(w), 640(sh), 620(s), 510(vw), 485(vw), 375(s), 358(vw), 350(vw), 310(s), 305(sh), 280(sh), 240(sh).
- $Bi_2O_3$ : no ir bands above 600 cm<sup>-1</sup>; below 600 cm<sup>-1</sup> broad bands on strong background bands at 545, 510, 450(sh), 430, 340, and 270.
- BiO(NO<sub>3</sub>): band in ir at 1380 (s), 835, and 820(w) presumably NO<sub>3</sub><sup>-</sup> ( $\nu_3$ ,  $\nu_2$ , and  $\nu_4$ ), 540(s,b), 380(s,b), and 275(s,b).



FIG. 6. Infrared and Raman spectra for Bi-molybdates; (a)  $Bi_2MoO_6(\gamma)$  (b)  $Bi_2MoO_6(\gamma')$ , (c)  $Bi_2Mo_2O_9(\beta)$ , and (d)  $Bi_2Mo_3O_{12}(\alpha)$ .

### DISCUSSION

Interpretation of ir and raman spectra. The crystal structures of the bismuth molybdates have been studied with precision by Van den Elzen and Rieck (16-18). Two structures, viz., those of Bi<sub>2</sub>MoO<sub>6</sub> ( $\gamma$ ) and Bi<sub>2</sub>Mo<sub>3</sub>O<sub>12</sub> ( $\alpha$ ), are now known with considerable accuracy from X-ray diffraction analysis. An "outline model" was given for

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$Bi_2MoO_6(\gamma)$		$\frac{\text{Bi}_2\text{MoO}_6 + 4\% \text{ MoO}_3}{(\text{extra bands})}$		$Bi_2MoO_6(\gamma')$		$Bi_2Mo_2O_9(\beta)$		Bi₂Mo	<sub>3</sub> Ο <sub>12</sub> (α)
ir	Raman	ir	Raman	ir	Raman	ir	Raman	ir	Raman
842 <sub>m</sub>	843 <sub>30</sub>		88616	890 <sub>sh</sub>	892100	890 <sub>m</sub>	892100	948 <sub>m</sub>	955 <sub>10</sub>
798 <sub>w</sub>	799 <sub>100</sub>	832 <sub>sh</sub>		880 <sub>sh</sub>	883 <sub>30</sub>	840 <sub>sh</sub>	845 <sub>30</sub>	931 <sub>m</sub>	919 <sub>30</sub>
735 <sub>vs</sub>	715 <sub>15</sub>		809 <sub>sh</sub>	868 <sub>vs</sub>	86441	800 <sub>s</sub>	793 <sub>60</sub>	902 <sub>m</sub>	897 <sub>82</sub>
600 <sub>w</sub>		760 <sub>sh</sub>	782 <sub>sh</sub>		851 <sub>23</sub>	770 <sub>s</sub>		860 <sub>s</sub>	853 <sub>20</sub>
		710 <sub>w</sub>		825 <sub>m</sub>	82237	740 <sub>m</sub>	71510	847 <sub>s</sub>	836 <sub>18</sub>
				810 <sub>vs</sub>		710 <sub>sh</sub>		830 <sub>m</sub>	$811_{28}$
				$780_{w}$	787 <sub>25</sub>			715 <sub>m</sub>	
					77029			670 <sub>s</sub>	654 <sub>12</sub>
				740 <sub>m</sub>					
				710 <sub>s</sub>	695 <sub>25</sub>				
				645 <sub>m</sub>					
570 <sub>m</sub>		545 shifted		562m		500 <sub>sh</sub>		570 <sub>m</sub>	
•••		from 570		527 w				550 <sub>sb</sub>	
450 <sub>m</sub>				467 <sub>w</sub>		473 <sub>m</sub>		470 <sub>m</sub>	
						446 <sub>m</sub>		440 <sub>m</sub>	
								415 <sub>sh</sub>	
408 <sub>m</sub>		$408_{st}$		391 <sub>w</sub>	399 <sub>9</sub>	409 <sub>m</sub>		392 <sub>m</sub>	
390 <sub>sh</sub>	399 <sub>5</sub>	390 <sub>st</sub>		380 <sub>w</sub>	37513	392 <sub>w</sub>		350 <sub>w</sub>	36018
352	354,	358 <sub>st</sub>		335 <sub>w</sub>	35910	362 <sub>m</sub>	35615	314 <sub>m</sub>	317,
304 <sub>m</sub>	32615			300 <sub>w</sub>	31525	$325_{vw}$	32316	$303_{vw}$	302 <sub>3</sub>
287 w	295 <sub>13</sub>			$290_{vw}$	303 <sub>27</sub>	275 <sub>m</sub>	286 <sub>21</sub>	$276_{vw}$	282 <sub>8</sub>
277 <sub>vw</sub>	27818				25816	266 <sub>w</sub>	260 <sub>5</sub>	$262_{vw}$	2485
268 <sub>vw</sub>					226 <sub>9</sub>		1802		
248 <sub>vw</sub>					206,		-		
	194 <sub>5</sub>								
	0								

### TABLE 2

## Infrared and Raman Spectra of Bismuth Molybdates

## Crystallographic Data for Bismuth Molybdates

	$Bi_2MoO_6$ Koechlinite ( $\gamma$ )	Blasse (y')	$\operatorname{Bi}_{2}\operatorname{Mo}_{2}\operatorname{O}_{9}(\beta)$	$\mathrm{Bi}_{2}\mathrm{Mo}_{3}\mathrm{O}_{12}\left(lpha ight)$
Crystal system	Orthorhombic	Orthorhombic	Orthorhombic	Orthorhombic
Space group	$Pca 2_1$	$P 2_{1}/c$	$P 2_1/c$	$P 2_1/c$
Coordination of Mo <sup>6+</sup>	Distorted octahedral	Tetrahedral	Trigonal- bipyramidal or tetrahedral	Trigonal- bipyramidal or distorted tetrahedral
Mo-O	$Bi_2MoO_6(\gamma)$	$Bi_2MO_3O_{12}(\alpha)$		
bond length (Å)	γ-	$\alpha_1$ -	$\alpha_{2}$ -	α3-
	Mo-O <sub>1</sub> 1.86	Mo <sub>1</sub> -O <sub>4</sub> 1.69	Mo <sub>2</sub> -O <sub>1</sub> 1.72	Mo <sub>3</sub> -O <sub>12</sub> 1.68
	-O <sub>4</sub> 2.24	$-O_5$ 1.72	–O <sub>9</sub> 1.73	-O <sub>11</sub> 1.78
	-O <sub>4</sub> 1.76	$-O_2$ 1.85	-O <sub>6</sub> 1.86	-O <sub>8</sub> 1.86
	-O <sub>5</sub> 2.24	$-O_{10}$ 1.92	-O <sub>3</sub> 1.89	-O <sub>7</sub> 1.87
	$-O_5 1.75$ $-O_6 1.93$	-O <sub>10</sub> 2.25	-O <sub>8</sub> 2.13	-O <sub>6</sub> 2.30



FIG. 7. Infrared and Raman spectra for "impregnated" samples; (a) stoichiometric  $Bi_2MoO_6(\gamma)$ , (b) 1 atom% of excess  $MoO_3$ , (c) 2 atom%, and (d) 4 atom%. Dashed line indicates  $Bi_2Mo_2O_9(\beta)$ .

 $Bi_2Mo_2O_9$  ( $\beta$ ) for which only powder diagrams could be obtained. Crystallographic data are given in Table 3.

A common feature of all these structures is the existence of  $MoO_n$  and  $BiO_m$  polygons with several, but not necessarily all, oxygen shared. An extensive literature concerns the spectra on  $MoO_n$ , with *n* varying between 4 and 6. It is usually assumed that n = 6 (octahedral) for Bi<sub>2</sub>MoO<sub>6</sub>( $\gamma$ ) and n =5 (trigonal bipyramidal) for  $Bi_2Mo_3O_{12}$  ( $\alpha$ ). However, the inspection of the structures reveals that, for both cases, there are two "short" bonds and two "intermediate" bonds, with in addition, for compound  $\gamma$ , two "long" bonds, and for compound  $\alpha$ , one "long" bond. The long bonds in both cases refer to an oxygen bonded with a "short" bond to a neighboring Mo (see Table 3). It therefore appears appropriate to assume that all Mo-oxygen polygons encountered here are derived from tetrahedral MoO<sub>4</sub>. The particular type of distortion can then be considered characteristic for the crystal structure. The structure of the "surface layer", of different form from that of the bulk, may, if circumstances permit, be guessed from the  $MoO_4$  spectra. A more extensive analysis of the spectra from compounds with known structure is thus necessary.

A tetrahedral group  $(T_d)$  has four vibrational modes,  $\nu_1(A)$ ,  $\nu_2(E)$ ,  $\nu_3(F_2)$ , and  $\nu_4(F_2)$ . All are Raman active, while only the triply degenerative modes are ir active. A complete lifting of the degeneracies would give a system of nine bands. Since  $\nu_1(A)$  and  $\nu_3(F_2)$  are intrinsically connected with stretching and  $\nu_2(E)$  and  $\nu_4(F_2)$  with bending modes, it may be expected that every Mo-O<sub>4</sub> configuration produces two groups of bands, positioned at different parts of the spectrum.

Busey and Keller (19) give information about the position of the two groups. Powdered Na<sub>2</sub>MoO<sub>4</sub> has four Raman bands at 892(vs), 808(s), 381(m), and 303(m) cm<sup>-1</sup>, while there are two ir bands at 838 and 325 cm<sup>-1</sup>. This indicates that the stretching modes are around 800 cm<sup>-1</sup> and the bending modes between 300 and 400  $cm^{-1}$ . Only the former are therefore useful to us, since the bending mode range for our system is probably overlapped by Bi-O bands. Somewhat closer to our system is PbMoO<sub>4</sub> with the scheelite structure, and all Mo-O bands equal to 1.772 Å (20). Tarte and Liegeois-Duyckaerts (21) found four ir bands (786, 374, 307, 242 cm<sup>-1</sup>) and six Raman bands (869, 766, 743, 357, 350, and 317 cm<sup>-1</sup>). It is noteworthy that, contrary to Bi-molybdates, the 400 to  $600\text{-cm}^{-1}$  range is empty, in good agreement with our assumption that bands in this range belong to Bi-O vibrations. It seems clear that there is sufficient evidence to accept that bands around 550  $cm^{-1}$  belong to this category. These bands are observable in BiO(NO<sub>3</sub>) and in BiOX with X being Cl, Br, or I (22). Also, Hulett and Thornton (23) investigated Bi troponolate complexes in which the troponolate anion (Tr) is a bidendate ligand. They found for BiTr<sub>2</sub>Cl bands at 514 and 580 cm<sup>-1</sup> while BiTr<sub>3</sub> has a band at 570 cm<sup>-1</sup>.

According to the X-ray structure, compound  $\gamma$  has only one type of MoO<sub>4</sub> unit ( $\gamma$ ) and there should be only four bands around 800 cm<sup>-1</sup>. This is what is observed, although it is somewhat difficult to locate the position of the fourth band that could be either the very weak 600 cm<sup>-1</sup> or the somewhat stronger 570 cm<sup>-1</sup> band. In agreement with the presence of three different MoO<sub>4</sub> groups ( $\alpha_1$ ,  $\alpha_2$ , and  $\alpha_3$ ), the higher range of the compound  $\alpha$  spectra are very complicated with at least nine bands. A more extensive analysis is therefore necessary before the spectra can be completely explained.

In order to do so, we used the G-F matrix method. Atomic coordinates and bond lengths were derived from the X-ray structural analysis of Van den Elzen and Rieck (16, 17) (see Table 3). Stretching force constants were derived from the Cotton–Wing relation (24) between Mo–O distance and stretching force constants. Stretching– stretching interaction constants, bending force constants, and bending–bending interaction constants were taken equal to those of tetrahedral MoO<sub>4</sub> (0.55, 0.35, and 0.02 m dyn/Å (25)). The computer program used was written by Bril and Vogel (26, 27) and is an extension of the Schachtschneider NCA program. Initial results were promising enough to warrant further refinement by slightly modifying the Cotton-Wing relation. The final results are given in Table 4, where they are confronted with the experimental results. There is a satisfactory agreement between predicted and observed band positions. Also, they clearly show that the range 600-400  $\text{cm}^{-1}$  is now empty, the bending modes being situated in the range 375–200 cm<sup>-1</sup>. The various experimentally observed bands can now be properly assigned to the various MoO<sub>4</sub> groups, as indicated in Table 4.

The agreement between prediction based on the tetrahedral model and experiment shows conclusively that the  $MoO_4$  model is entirely sufficient to describe the spectra. It is therefore doubtful whether evidence from ir and Raman spectroscopy can be used to make a distinction between "octahedral" or "tetrahedral" surrounding of the molybdenum. It does, however, allow the recognition of a greater or lesser degree

TABLE 4

Calculated Band Positions for  $MoO_4$  of  $Bi_2MoO_6(\gamma)$  and  $Bi_2Mo_3O_{12}(\alpha)$  and Comparisons with Observed Bands (cm<sup>-1</sup>)

Assignment	$\frac{\text{Bi}_2\text{Mo}_3\text{O}_{12}(\alpha)}{\text{MoO}_4(\alpha_1) \text{ MoO}_4(\alpha_2) \text{ MoO}_4(\alpha_3)}$					$O_6(\gamma)$	) of Bi <sub>2</sub> MoC	<b>ΜοΟ</b> 4 (γ
								Calculated
	Observed					served	Obs	
	Raman	ir	đ	Calculated		Raman	ir	
	955	948	933			848	842	865
$\alpha_1$	919	931			921	799	798	780
$\alpha_2$	897	906		886		715	735	693
$\alpha_{i}$	853	860			849		600	613
$\alpha_2$	836	847		822				
$\alpha_3$	811	830	796					
$\alpha_{1,2,3}$		715	690	690	703			
$\alpha_{1,2,3}$	654	670	680	665	651			
	360	350	341	343	345	399	375	376
	317	314	334	340	336	354	354	355
	302	303	330	336	335	326	304	337
	282	276	297	292	295	293	287	287
	248	262	263	249	295	194		1 <b>9</b> 6

of distortion of the pseudotetrahedral configuration. For instance, taking the accepted structure of  $\gamma$  and  $\alpha$ , as determined by X-ray diffraction, one can predict that the former should have a Mo-O-Bi mode of vibration near 600 cm<sup>-1</sup> but not the latter. Since both have a 570-cm<sup>-1</sup> band but only the former has a band (very weak), only this latter band should be attributed to the MoO<sub>4</sub> stretching mode. The quantitative agreement (theory, 613; experiment, 600 cm<sup>-1</sup>) confirms this assignment.

Let us now discuss the spectra of Bimolybdates for which the structure is only approximately known ( $\gamma'$ ,  $\beta$ ). For the  $\gamma'$ -Bi<sub>2</sub>MoO<sub>6</sub>, Blasse (28) considers its structure to be similar to that of koechlinite  $(\gamma)$ except for the arrangement in the MoO<sub>2</sub>layer. He suggests this to consist of separate MoO<sub>4</sub>-groups. The complexity of the spectra of  $\gamma'$  indicates the presence of at least two, but possibly three, different  $MoO_4$  configurations. The prediction of Blasse (28) that the structure is similar to that of koechlinite  $(\gamma)$  is therefore not entirely acceptable. In fact, the  $\gamma'$ -spectrum is more similar, although not identical, to that of  $Bi_2Mo_3O_{12}(\alpha)$ .

Van den Elzen and Rieck (18), in their outline of a model for the  $Bi_2Mo_2O_9$  ( $\beta$ ) structure, consider all MoO<sub>4</sub>-units as equivalent; we showed above that there might be two different types of MoO<sub>4</sub>, although the extent to which they are different cannot be specified. Supposing, according to their model, that compound  $\beta$  contains building elements such as are present in compound  $\gamma$  and  $\alpha$ , then the two MoO<sub>4</sub> structures might resemble  $\gamma$ -MoO<sub>4</sub> and one of the  $\alpha$ -MoO<sub>4</sub> units. Indeed, three of the four  $\gamma$ - $MoO_4$  bands (840, 790, and 740 cm<sup>-1</sup>) are observable in the spectra of  $\beta$ , while the fourth band at 600  $cm^{-1}$  is very weak. So we propose that there is a  $\gamma$ -set similar to the  $\beta_1$ -set. It is also remarkable that not only the position but also the ir and Raman intensities of the genuine  $\gamma$ -set and that assumed in the  $\beta_1$ -set agree, the 790 to 800cm<sup>-1</sup> band being strong in Raman but weak

in ir while the 735-cm<sup>-1</sup> band is strong in ir with no Raman band observable at this point. The remaining bands, 890(R), 770(ir), and 715(ir) with 840-cm<sup>-1</sup> (R, ir) band presumed to make up the  $\beta_2$ -set, have some similarity with the  $\alpha_2$ -set in Bi<sub>2</sub>Mo<sub>3</sub>O<sub>12</sub> ( $\alpha$ ) but the 770-cm<sup>-1</sup> band is absent in the compound  $\alpha$ .

Our proposal is similar to that of Gryzbowska *et al.* (29) and of Mitchell and Trifirò (30), except for the fact that they consider  $Bi_2Mo_2O_9$  ( $\beta$ ) to have tetrahedral and octahedral Mo which we choose to be more specific by insisting that the tetrahedral surrounding is actually a distorted trigonal bipyramidal arrangement ( $\alpha_2$ -set) and "octahedral" Mo is strongly distorted, similar to that of the  $\gamma$ -set.

We now compare the positions of incipient bands formed by the introduction of small amounts of MoO<sub>3</sub> in Bi<sub>2</sub>MoO<sub>6</sub> ( $\gamma$ ) with those of the bands of the compounds. We start with the set of extra bands in the range above 600  $cm^{-1}$  (see Table 2 and Fig. 7, where the ir and Raman bands of Bi<sub>2</sub>Mo<sub>2</sub>O<sub>9</sub>  $(\beta)$  are confronted with those of the impregnated samples). It is obvious that there is no correspondence with the band positions for the  $MoO_3$ -set. There is, however, a distinct similarity with the  $\beta_2$ -set. Two strong bands for this compound, the strong 892-cm<sup>-1</sup> Raman band and the massive 770cm<sup>-1</sup> ir band, coincide with a tiny Raman peak at 886 cm<sup>-1</sup> and a rather strong shoulder around 760 cm<sup>-1</sup> (ir) on the  $\gamma$ -band. A small but distinct shoulder at  $832 \text{ cm}^{-1}$  (ir) in the extra set was not traced in the  $\beta$ spectrum, but could have been overlooked because it fused with the 840-cm<sup>-1</sup> band ( $\gamma$ =  $\beta_1$ ). If so, the complete  $\beta_2$ -set would be 890-892, 832, 770, 710-715, and 473 cm<sup>-1</sup>. The extra set could not be derived from  $Bi_2Mo_3O_{12}$ , because the  $\alpha_1$  and  $\alpha_3$  bands are completely lacking and  $\alpha_2$ , although slightly resembling  $\beta_2$ , is definitely different as remarked before. We therefore conclude that the incipient set of bands is that of a surface layer with a structure closely resembling that of  $Bi_2Mo_2O_9$ ; the set is very similar to

 $\beta_2$  while  $\beta_1 = \gamma$  remains hidden by the more massive Bi<sub>2</sub>MoO<sub>6</sub> spectrum.

Next we turn to the change occurring in the  $Bi_2MoO_6$  spectra below 600 cm<sup>-1</sup> (see Fig. 6a and Table 2). These spectra are different in character from those considered above. Above  $600 \text{ cm}^{-1}$ , we find the formation of new bands of low intensity next to the more massive bands of the mother compound. We also observe bands already present in the Bi<sub>2</sub>MoO<sub>6</sub> spectra that shift positions and become more intense. The general impression is that of an increase in regularity. While the new bands above 600  $cm^{-1}$  could be assigned to surface layers, the changes below 600 cm<sup>-1</sup> are undoubtedly connected with changes in the bulk, certainly an unusual phenomenon considering the small amounts of MoO<sub>3</sub> incorporated. Even more surprising is that incorporation of more substantial amounts, up to 30%, does not lead to further changes.

The changes occur in two regions, viz.:

(i) Betwen 570 and 500 cm<sup>-1</sup>, i.e., where the Bi-O bands are supposedly situated. The weak 570-cm<sup>-1</sup> band moves to 545 cm<sup>-1</sup> and becomes more intense. Hulett and Thornton (23) observe similar shifts in passing from BiTr<sub>2</sub>+Cl<sup>-</sup> to BiTr<sub>3</sub> (from 514 to 495 cm<sup>-1</sup>). They ascribe this to an increase in coordination number from 4 to 6.

(ii) In the Mo-O bending vibration range at 400-200 cm<sup>-1</sup>, a rather broad and irregular band in the ir below 410 cm<sup>-1</sup> grows to a system of three equally intense bands at 410, 390, and 358 cm<sup>-1</sup>. Similarly, in the Raman we now find three quite distinct bands at 354, 326, and 295 cm<sup>-1</sup>.

The following tentative explanation is offered to account for the phenomena. We start from the assumption that samples of stoichiometric  $Bi_2MoO_6$  contain defects of a type in which Mo and Bi are in the wrong cation positions. Interdiffusion of cations would eliminate the defects, but this could be supposed a difficult process, since there are no cation vacancies. However, once vacancies are produced by the incorporation of small amounts of Mo, the defect lattice can relax, because cations can now migrate. In other words, the spectra produced after Mo incorporation are those that actually belong to the defect-free Bi<sub>2</sub>MoO<sub>6</sub>. The ir group might be related to the  $\nu_4(F_{1u})$ bending mode of O<sub>h</sub> (only infrared active) and the Raman group to  $\nu_5(F_{2u})$ .

Introduction of cation vacancies by Mo incorporation can be supposed to "catalyze" other phase transitions. We found indeed that the koechlinite to Blasse phase transition occurs at temperatures appreciably lower for  $Bi_2MoO_6$  + small amounts of  $MoO_3$  (580°C) than for stoichiometric  $Bi_2MoO_6$  (640°C).

Surface structure and catalytic activity. The combined evidence of XPS and ir and Raman spectroscopy strongly supports the hypothesis that the catalytic activity of Bimolybdate system resides in "surface domains" that may be different in structure and composition from the bulk. Here we postulate that the "surface domains" are similar in structure and composition to  $Bi_2Mo_2O_9(\beta)$ . The following arguments underline this hypothesis:

(1) While compound  $\gamma$  (stoichiometric) is almost inactive and compound  $\alpha$  has low activity, all mixtures of the two in the range 2/3 < Bi/Mo < 2/1 are very active and have an activity similar to that of compound  $\beta$ .

(2) Even relatively small amounts of  $MoO_3$  in excess of the stoichiometric ratio suffice to bring the catalyst to maximal activity:  $Bi_2Mo_{1.04}O_{6.12}$  is just as active as compound  $\beta$ .

(3) XPS data show the existence of a surface layer with a composition different from that of the bulk phase, the Bi/Mo ratio at the surface in the high-activity range being almost constant at a level near 1.

(4) The ir and Raman spectra of catalysts such as  $Bi_2Mo_{1.04}O_{6.12}$  show a new, weak set of bands in the stretching mode region of Mo-O vibrations. This set is different from that of pure MoO<sub>3</sub>; this shows that superficial MoO<sub>3</sub> has reacted with the underlying  $Bi_2MOO_6(\gamma)$ . This set of bands is similar to one of the sets  $(\beta_2)$  that can be observed in the spectrum of  $\beta$ .

(5) Since the other band system of this compound  $(\beta_1)$  is akin to that of compound  $\gamma$ , a layer of  $\beta$  on  $\gamma$  would only show the  $\beta_2$ -set superposed on the  $\gamma$ -set, as was observed.

For the structure of  $Bi_2Mo_2O_9$  ( $\beta$ ), the model (18) proposed is a combination of elements derived from the Scheelite and from the koechlinite  $(\gamma)$  structure. The Scheelite-derived contribution is a stacking of four Mo, in a square pattern, bound together by Bi positioned on the axis passing through the centers of the squares and half-way between two groups of four Mo ("Mo<sub>4</sub>"). It is assumed that the "Mo<sub>4</sub>" arrangement represents Mo<sub>4</sub>O<sub>16</sub> clusters bonded together by bridging oxygens, as in the Mo<sub>2</sub>O<sub>8</sub> groups in compound  $\alpha$ . The particular Bi-cations (Bi<sup>1</sup>) are therefore surrounded by eight  $MoO_4$  units, as is  $Pb^{2+}$  in  $PbMoO_4$ , the difference being that the distance between the MoO<sub>4</sub> units is smaller in two crystallographic directions (x and z). A typical koechlinite  $(\gamma)$  building element is given by the linear rows of oxygens, running parallel to the  $(Mo_4O_{16})$ -Bi- $(Mo_4O_{16})$ axis. These oxygens are only connected to Bi-cations and are similar to ribbons of Bi<sub>2</sub><sup>11</sup>O cut out from the Bi<sub>2</sub>O<sub>2</sub>-planes in koechlinite, while the Bi<sup>II</sup>-cations are further linked to the oxygens of the Mo<sub>4</sub>O<sub>16</sub>groups. A noteworthy detail is that not all these Bi-sites are filled, one in every four sites being empty. The model might be summarized by writing it as

$$Bi^{I}(Bi_{3}^{II}\phi_{c}O_{2})(Mo_{4}O_{16}),$$

 $\phi_{\rm c}$  being the empty cation site.

As we discussed earlier, the structure of  $\beta$  thus appears more closely related to the structure of  $\gamma$  than believed so far. Figures 8a and b suggest a method to illustrate the structural relation. Figure 8b gives cation and cation vacancy positions in the (010) plane of the  $\beta$  structure (18). The crystal planes (101) perpendicular to this cation plane can now be related to the (Bi<sub>2</sub>O<sub>2</sub>) and

 $(MoO_2)$  planes in compound  $\gamma$  as shown in Fig. 8a. To go from Fig. 8a to b, pairs of Bications have to be replaced by one Mo and one cation vacancy, but there is simultaneously some reshuffling of cations according to the scheme

$$(Bi_2)(Mo) \rightarrow (BiMo_{2/3}\phi_{1/3})(Mo_{2/3}Bi_{1/3}),$$

i.e., Bi-cations are replaced by Mo in the Bi<sub>2</sub>O<sub>2</sub> layer but simultaneously some Mo has been replaced by Bi in the MoO<sub>2</sub> layer. If we assume the oxygen positions to remain similar to those in  $Bi_2MoO_6(\gamma)$  we arrive at Fig. 8c. It is then obvious that there are two different oxygen surroundings of Mo, one equal to the  $\gamma$ -surrounding, the other a new one in the Bi<sub>2</sub>O<sub>2</sub>-derived layer. Of course, since the  $\gamma$ -structure chosen as position represents the starting а simplification of the more distorted real structure, it is probable that similar distortions will occur in the  $\beta$ -structure. However, the model is adequate to expalin the spectral observation as to its essential points.

A sample with a bulk composition Bi/Mo = 1.96 has sufficient excess MoO<sub>3</sub> to cover the entire Bi<sub>2</sub>MoO<sub>6</sub> surface with a layer of MoO<sub>3</sub>. XPS data show, however, that the surface Bi/Mo ratio remains as high as 1.5. Noble-gas ion reflexion mass spectrometry (NIRMS) studies carried out by courtesy of Professor H. H. Brongersma also indicate an almost equal ratio (Bi/Mo = 1.6). According to calculations based on the experi-



F1G. 8. Cation positions for (a)  $Bi_2MoO_6(101)$  plane, (b)  $Bi_2Mo_2O_9$  (010) plane, and (c) modified  $Bi_2Mo_2O_9$ (010) plane with oxygen positions.

mental surface area and density, the thickness of the  $\beta$ -phase overlayer at 4% excess Mo on Bi<sub>2</sub>MoO<sub>6</sub> should be about 10 Å. Since NIRMS is confined to the outer layer (penetration depth < 5 Å), while XPS probes deeper (15 ~ 40 Å), this indicates that the catalyst seems to form either a thicker  $\beta$ -phase island on Bi<sub>2</sub>MoO<sub>6</sub> or that the surface layer corresponds to a ratio of 1.5. In the latter case, the surface structure would be

$$\begin{array}{c} Bi_2MoO_6 \xrightarrow{+Mo} \\ (Bi_{1.57}Mo_{0.287}\phi_{0.143})(Mo_{0.857}Bi_{0.143})O_6. \end{array}$$

The surface can be considered as a  $Bi_2MoO_6$ -surface with a number of "surface defects" such as Mo-cations and cation vacancies in the  $Bi_2O_2$ -layer and Bi ions in the  $MoO_2$ -layers. The defects may be either randomly distributed over the surface or ordered in  $Bi_2Mo_2O_8$ -like "surface domains". Here, however, we do not touch on which surface state is actually formed. The answer to this question would necessitate further investigations using, for example, XPS with ion etching or electron microscopic studies.

These data may be used to construct an "outline surface model". Indeed, one of the authors (31) had earlier proposed a model based on the assumption that  $Bi_2MOO_6$  is an active catalyst. The surface was assumed to consist entirely of planes (010) perpendicular to the  $MOO_{2^-}$  and

 $Bi_2O_2$ -layers. We have given the proof that this particular plane in stoichiometric  $Bi_2MoO_6(\gamma)$  cannot be active since the compound is not active at all and only becomes so by introducing extra Mo. Consequently, our active surface plane has to be derived from the (010) plane of  $Bi_2Mo_2O_3(\beta)$ . Figure 9 shows the formation of "surface domains  $(\beta)$ " supposed to be the active species, at the (101) plane of  $\gamma$  when excess Mo is introduced in  $\gamma$ . The process of removing oxygens from the top layer (Fig. 8c) has been performed by discarding one out of every two equivalent oxygens. The surface model shows that the olefin adsorption site (B-site) in the MoO<sub>2</sub>-layer discussed earlier (31) still exists in the "surface domains". It has been shown that the adsorption parameters for but-1-ene both on stoichiometric and Mo-enriched active  $\gamma$  are similar (see Table 1). Also, we assumed an active oxygen  $(O_A)$  should be located at the edge of the  $Bi_2O_2$  plane in  $Bi_2MoO_6(\gamma)$  (31). For the  $Bi_2O_2$  layer we see that the layer is formed from alternating vacancies and oxygens, each vacancy and oxygen belonging to a pair of Bi cations. Some oxygens of the "surface domains" are only connected with one Bi-cation; the other site is empty. This type of oxygen should be active and form  $O_A$ . Correlation between activity and the maximum volume of strong butadiene adsorption  $(V_m)$  which relates to  $O_A$  is shown in Fig. 10. It seems clear that the formation of cation vacancies in



FIG. 9. Surface model of "surface domains" at the Bi<sub>2</sub>MoO<sub>6</sub> (101) plane.



FIG. 10. Dependence of catalytic activity on the maximum volumes of strong butadiene adsorption  $(V_m)$  related to  $O_A$  on "impregnated" samples.

 $Bi_2MoO_6(\gamma)$  with addition of Mo makes the catalyst active.

The postulate that only  $Bi_2Mo_2O_4(\beta)$  is the active component agrees with the theory of Sleight and Linn (32) that a necessary feature of an active catalyst is the presence of cation vacancies in sites normally occupied by  $Bi^{3+}$  as given in the molecular formula  $Bi^1$  ( $Bi_3^{II}\phi_c^{II}O_2$ )( $Mo_4O_{16}$ ). If  $Bi^1$  is allowed to move over to the cation vacancy, we obtain  $\phi_c(Bi_4^{II}O_2)(Mo_4O_{16})$ . The I position would result in a situation entirely similar to that postulated by Sleight and Linn (32) to occur when PbMoO<sub>4</sub> is doped with  $Bi^{3+}$ . Another possibility for randomly occurring vacancies can be frequent twinning of  $Bi_2Mo_2O_9$  crystals.

Our above conclusions rest on a combination of measurements concerning activity, adsorption, XPS and optical spectroscopic information. They are only strictly valid for samples very near the Bi/Mo = 2edge for Bi<sub>2</sub>MoO<sub>6</sub>. For the majority of the "multiphase" samples in the entire range  $2/3 \le \text{Bi}/\text{Mo} \le 2/1$ , we have only the XPS data, the Bi/Mo  $\approx$  1 value being indicative of the presence of a surface layer of the corresponding composition. Spectroscopic information ceases to be applicable because of the simultaneous presence of the  $\gamma$ - and  $\alpha$ -phases, causing the ir and Raman spectra to become so complicated that identification of a separate surface layer becomes nearly impossible. Therefore, the presence of another type of surface layer with a composition  $Bi/Mo \approx 1$  cannot be entirely excluded in the  $2/3 \leq Bi/Mo \leq 2/1$ range, although it is believed to be improbable.

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